Double-peaked Oxygen Lines Are not Rare in Nebular Spectra of Core-Collapse Supernovae

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ABSTRACT

Double-peaked oxygen lines in the nebular spectra of two peculiar Type Ib/c Supernovae (SN Ib/c) have been interpreted as off-axis GRB-jet or unipolar blob ejections. Here we present late-time spectra of 8 SN IIb, Ib and Ic and show that this phenomenon is common and probably should not be linked to extraordinary events in the explosion physics. We show that this type of line profile is probably not caused by optical depth effects, but might well be due to ejecta expanding with a torus- or disk-like geometry. Double-peaked oxygen profiles are not necessarily the indicator of a mis-directed GRB jet.

Subject headings: gamma-ray burst; general supernovae: general

1. INTRODUCTION

To gain a full understanding of the connection between long-duration Gamma Ray Bursts (GRBs) and the supernovae (SN) with which some of them are known to occur, it would be significant to detect the presence of jets produced in the explosion when the observer is not in the beam (for a review, see Woosley & Bloom 2006). In the broad-lined Type Ic, SN 2003jd, an off-axis GRB jet was proposed as the cause for the double-peaked oxygen profile observed in its nebular spectrum (Mazzali et al. 2005; Valenti et al. 2007). A similar double-peaked oxygen profile, with a large blueshift, was detected in the peculiar Type Ib, SN 2005bf (Maeda et al. 2007; Modjaz 2007), where it was interpreted as coming from a unipolar blob or jet (Maeda et al. 2007).

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If the double-peaked oxygen shape were exclusively caused by highly relativistic jets, then we might expect only broad-lined SN Ic, whose line widths approach 30,000 km s\(^{-1}\), and which are the only subtype of stripped-envelope SN seen at the sites of GRBs to exhibit them. However, if we see double-peaked profiles in a range of supernova types, we might conclude that asphericities are present in more typical core-collapse events.

We set out to obtain a set of nebular spectra for supernovae of various subtypes to measure the emission line profiles with adequate resolution and signal to detect signs of these double-peaked profiles in nebular spectra of supernovae. As the supernova turns optically thin a few months after maximum light, the emission line shapes can provide information on the velocity distribution of the ejecta (Matheson et al. 2001; Foley et al. 2003; Maeda et al. 2006). Structure in emission line profiles have been reported in SN II (Spyromilio 1991), SN IIb (Spyromilio 1994; Matheson et al. 2000) and SN Ib (Sollerman et al. 1998; Elmhamdi et al. 2004), and interpreted as clumping of oxygen in the ejecta.

Here we report on late-time spectra of 8 supernovae, of types IIb, Ib, and Ic, of which three show double-peaked profiles. In addition, we draw attention to double-peaked oxygen profiles in spectra of the classical SN Ib 1984L (Schlegel & Kirshner 1989) that have not previously been reported. In § 2 we discuss our observations and present our sample. In § 3 we discuss the line profiles in detail and investigate whether the double-peaked profiles could be caused by optical depth effects. We discuss and speculate on the implications of our results in § 4 and summarize in § 5.

2. Observations and Analysis

Spectra were obtained with the 6.5 m Clay Telescope of the Magellan Observatory located at Las Campanas Observatory (LCO), with Gemini-North via queue-scheduled observations (GN-2005B-Q-11, GN-2006B-Q-16 PI: Modjaz), with the 6.5 m Multiple Mirror Telescope (MMT) and the 1.5 m Tillinghast telescope at the Fred Lawrence Whipple Observatory (FLWO). The spectrographs utilized were the LDSS-3 (Mulchaey & Gladders 2005) at LCO, the GMOS-North (Hook et al. 2003) at Gemini, the Blue Channel (Schmidt et al. 1989) at the MMT, and FAST (Fabricant et al. 1998) at the FLWO 1.5 m telescope. All optical spectra were reduced and calibrated employing standard techniques in IRAF\(^5\) and our own IDL routines for flux calibration (see e.g., Matheson et al. 2008). In Figure 1, we

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present the observed nebular spectra of a total 8 newe SN IIb, SN Ib and SN Ic. For clarity we show only the latest spectrum of each SN (except for SN 2004ao). Thus, we double the number of unambiguous SN Ib with late-time data from four cases (SNe 1983N, 1984L, 1990I, and 1996N) to eight, not counting SN 2005bf (see also Maeda et al. 2008). We determined the phase with respect to maximum from our own photometry (Modjaz 2007, M. Modjaz, in prep.). In addition, we show our data of the peculiar SN Ib 2005bf, independently observed and analyzed by Maeda et al. (2007), and of SN 2006aj, the broad-lined SN Ic connected with GRB 060218 (Modjaz et al. 2006), independently observed and analyzed by Mazzali et al. (2007).

All spectra display the hallmark features of nebular SN Ib/c spectra: strong forbidden emission lines of intermediate-mass elements such as [O I] $\lambda\lambda 6300,6363$ and [Ca II] $\lambda\lambda 7291,7324$, similar to spectra of SN II. In the absence of hydrogen, oxygen is the primary coolant in the ejecta of stripped-envelope SN at late epochs when the gas is neutral or at most singly ionized (Uomoto & Kirshner 1986; Fransson & Chevalier 1987). The strong emission line at 7300 Å could instead be due to [O II] $\lambda\lambda 7319,7330$, or be a blend of both [O II] and [Ca II] (see Fesen et al. 1999 and references therein). Nevertheless, we do not believe that [O II] is contributing significantly, as we would expect to otherwise see strong [O III] $\lambda 5007$, which was detected in SNe II 1979C and 1980K and powered by circumstellar interaction. SN 2005bf is the only SN discussed here that exhibits broad Hα (see also Maeda et al. 2007).

3. Detection of Double-peaked Oxygen Lines

In the optically-thin case, the late-time emission line profile is in principle dictated by the geometry and distribution of the emitting material (Fransson & Chevalier 1987; Fransson 1988; Schlegel & Kirshner 1989). A radially expanding spherical shell of gas produces a square-topped profile, while a filled uniform sphere produces a parabolic profile. In contrast, a cylindrical ring, or torus, that expands in the equatorial plane gives rise to a “double-horned” profile as there is very little low-velocity emission in the system, while the bulk of the emitting gas is located at ±$v_t$, where $v_t$ is the projected expansion velocity at the torus.

When zooming in on the specific line shapes in Figure 1, we note that certain SNe (SNe 2004ao, 2004gt, 2006T) show a conspicuous double-peaked line profile. For clarity, we plot the subset of SN with double-peaked oxygen lines in velocity space, with respect to 6300 Å in Figure 2. A literature search of published SN IIb, Ib, and Ic spectra reveals that SN Ib 1984L, the prototype of the SN Ib class, also shows a clear double-peaked profile that went unremarked in the original publication (Schlegel & Kirshner 1989). From the available three low-resolution late-time spectra of SN 1984L, we plot the one with the highest
S/N in Figure 2. Finally, we also include our data on SN 2005bf, the peculiar SN Ib with extraordinary early-time light curves and spectra (Tominaga et al. 2005; Folatelli et al. 2006), that showed highly blue-shifted (by \( \sim 2000 \text{ km s}^{-1} \)) oxygen and calcium lines (Maeda et al. 2007; Modjaz 2007).

This double-peaked line profile is visible in the oxygen line and not in calcium. For SN 2004ao, the only SN for which we have a sufficiently high S/N spectra, the double-horned feature is present in the permitted line of oxygen, \( \text{O I} \lambda 7774 \), and in \( \text{[Mg I} \lambda 4571 \) (see Fig. 2). Our late-time spectral follow-up for some SN (Modjaz 2007) shows that SN without double-peaked profiles (SNe 2004gk, 2004gq) do not exhibit them at any epoch, while SN with double peaks (SNe 2004ao, 2004gt) retain them over the full observed period.

The two horns are roughly symmetrically offset by \( \sim 1000\text{–}2000 \text{ km s}^{-1} \) around the trough. While for most SN in our sample the trough between the two “horns” is located at nearly zero velocity (with respect to 6300 Å), certain SN exhibit an apparent blueshift (SN 2004ao) or redshift (SN 1984L). We discuss possible causes for such shifts in § 4. We estimate that a velocity shift of up to \( \pm 200 \text{ km s}^{-1} \) (as seen in SN 2004gt) could simply be due to differences between the measured redshift of the nucleus of the host galaxy in NED (which was used to correct the SN spectra) and the actual radial velocity at the position of the SN, caused by rotation curves in their host spiral galaxies or line-of-sight coincidences.

### 3.1. Optical Depth Effects causing the Double Peaks?

Could the observed double-peaked profile be due to optical depth effects? The \( \text{[O I} \lambda \lambda 6300,6363 \) doublet has a velocity separation of 2998 km s\(^{-1}\) (with respect to 6300 Å) and an intensity ratio of 3:1 in the optically thin limit. Optical depth effects for \( \text{[O I} \lambda \lambda 6300,6363 \) have been discussed in the supernova context by Leibundgut et al. (1991), Spyromilio (1991) and Li & McCray (1992). In these models, the ratio of \( \text{[O I} \lambda 6300 \) to \( \text{[O I} \lambda 6363 \) evolves from \( \sim 1 \) to \( \sim 3 \) from early times to late times as the supernova expands and the lines become optically thin, as was observed for SNe 1987A (Spyromilio et al. 1991; Li & McCray 1992) and 1988A (Spyromilio 1991). In Figure 3, we plot the measured ratio of the blue to the red horn for SN 2004ao, the supernova in our sample with the longest time span of observations. The ratio of line intensity decreases over time for SN 2004ao. In contrast, this ratio would increase with time if the two horns were due to \( \text{[O I} \lambda 6300 \) and \( \text{[O I} \lambda 6363 \), as seen in SN 1987A (Li & McCray 1992, overplotted in Fig. 3).

In SN 2004ao, where we have sufficiently high S/N- data, we also see a similar double-peaked profile in the permitted \( \text{O I} \lambda 7774 \) line and in \( \text{[Mg I} \lambda 4571 \) which are single-transition
lines, as shown in Fig. 2. This suggests that the line shapes really do provide a guide to the distribution of oxygen and magnesium in the ejecta. Thus, we conclude that optical depth is probably not the cause of seeing two more-or-less equal peaks in the \([\text{O I}] \lambda\lambda 6300,6363\) line.

4. Interpretation

Our findings suggest a torus-like structure for one (or both) of the following distributions: either 1) the oxygen distribution or 2) the distribution of \(^{56}\text{Co}\), which is the main radioactive species for energy deposition and source of exciting the lines at these epochs (e.g., Fransson & Chevalier 1989). If 2) were the case, we would expect to see the same double-peaked profile in the calcium line, too, because the \(\gamma\)-rays produced during the decay of \(^{56}\text{Co}\) are powering both calcium and oxygen emission lines. Since the double-peaked profile is only observed in oxygen, we believe 2) is less likely, and we favor 1), namely that the bulk of the oxygen distribution is located at projected velocities of \(\sim 1000-2000\) km s\(^{-1}\). Models of the line shape that results from an expanding disk are worked out by Leonard et al. (2000), Gerardy et al. (2000), and Fransson et al. (2005). Comprehensive modeling of the line profiles needs to take into account viewing angle effects and the thickness of the line-emitting and expanding region. We leave that as a topic for future work.

Here we do not attempt to explain the extreme blueshift (\(\sim 2000\) km s\(^{-1}\)) seen in SN 2005bf at epochs later than 200 days. The less extreme blueshifts seen in spectra obtained at \(t \lesssim 200\) days for one of our SN (SN 2004ao) might be caused by the same mechanism that is responsible for the observed blueshifts in SN that show a single-peak line of \([\text{O I}] \lambda\lambda 6300,6363\). S. Taubenberger et al. (in preparation) find in a large set of 100 late-time spectra of 34 SN Ib and SN Ic that the oxygen line centroids are found to be blueshifted for spectra taken at \(t \lesssim 200\) days. These blueshifts range up to \(\sim 1500\) km s\(^{-1}\) for spectra at \(t \sim 90\) days and go to zero with increasing time. Taubenberger et al. exclude dust formation, contamination from other lines and geometric effects as potential causes and invoke residual opacity effects as the most likely reason. The observed blueshift-values in our sample agree with those seen in Taubenberger et al. for similar epochs; thus, the same mechanism might be causing the blueshifts in both single- and double-peaked \([\text{O I}] \lambda\lambda 6300,6363\). We encourage future detailed modeling of the radiative transfer of \([\text{O I}] \lambda\lambda 6300,6363\) to elucidate the exact reason for the blueshift. We have no simple explanation for the redshift seen in the trough of the double-peaked \([\text{O I}] \lambda\lambda 6300,6363\) of SN 1984L.

We conclude that the most probable explanation for the double-horned oxygen profiles is a torus like, or at least flattened, distribution of oxygen, that may be a relic of the explosion
physics. Energy deposition and radiation transport modeling of late-time emission in SN Ib (Fransson & Chevalier 1989) predicts that material contributing to [O I] λλ6300,6363 (and to [Mg I] λ4571) emission is situated further out in the ejecta than material emitting in [Ca II] λλ7291,7324 (see their Fig. 8). We speculate that the mechanism causing the asphericities has to affect the outer layers of the SN debris more than the inner ones. While Raleigh-Taylor or hydrodynamic instabilities are good candidates for causing anisotropies, they are expected to operate on small scales, i.e. induce small sub-structure superimposed on the oxygen lines, as seen in the SNe II 1979C, 1980K (Fesen et al. 1999 and references therein), in SN IIb 1993J (Matheson et al. 2000) and in SN Ib/c 1985F (Filippenko & Sargent 1986). In the case of our observed SN, the anisotropies may be of global nature in order to give rise to such a clear double-peaked line profiles. Alternatively, in case the supernova progenitors are part of binaries, binary interaction or merger might be modulating the geometry of the supernova ejecta (Morris & Podsiadlowski 2007).

Such a clear signature of a pure double-peaked oxygen profile has not been found in other SN Ib/c (Matheson et al. 2001), except in SN 2003jd (Mazzali et al. 2005) and in SN 2005bf. SN 2003jd was a broad-lined SN Ic without an observed GRB and showed the same double-peaked oxygen profile (Mazzali et al. 2005) as SN 2004ao (compare their Fig. 2 with our Fig. 2). Mazzali et al. (2005) interpret their observations as indicating an aspherical axisymmetric explosion viewed from near the equatorial plane. Moreover, they suggest that this asphericity was caused by an off-axis GRB jet. From computing nebular spectra using the same 2D explosion models of Maeda et al. (2006), as was done for GRB 980425/SN 1998bw, they find that the model with a jet directed > 70 deg away from our line-of-sight can reproduce the line profile well. Furthermore, they conclude that only special SN, those with high expansion velocities and possibly connected with GRBs, are aspherical, and use SN 2003jd as a link between the normal, spherical SN Ic and those highly aspherical ones connected with GRBs. In our sample, however, the SN showing these line profiles are normal SN Ib and even SN IIb, i.e. SN from stars with intact helium and partial H envelopes before explosion, and have normal early-time expansion velocities. Thus, it appears that asphericities are prevalent in normal core-collapse events.

Indeed, complex and ring-like velocity structures of oxygen have been observed in SN remnants (SNRs), most prominently in the oxygen-rich SNR 1E0102.2−7219 in the Small Magellanic Cloud (Tuohy & Dopita 1983). Its progenitor is suggested to have had oxygen-rich mantels, and possibly was a WR star (Blair et al. 2000). Tuohy & Dopita (1983) and subsequent papers fit a twisted ring model to the spatial and velocity extent of the filaments that extents to velocities ranging between −2500 to + 4000 km s$^{-1}$. It is conceivable that the stellar death leading to SNR 1E0102.2−7219 had a similar ejecta geometry found in our data set of SN. Thus, asphericities might be common in core-collapse events, be they neutrino-
driven (Scheck et al. 2006), acoustic (Burrows et al. 2006), or magneto driven (Burrows et al. 2007; Dessart et al. 2008), as indicated by multiple lines of evidence, such as polarization studies of SN II, Ib and SN Ic (Leonard & Filippenko 2005), neutron-star kick velocities (Wang et al. 2006) and young SN remnant morphologies (Fesen et al. 2006).

5. Conclusions

In summary, we detect clear double-peaked line profiles of [O I] \( \lambda \lambda 6300,6363 \) in three SN IIb and SN Ib/c out of our observed sample of 8 and additionally, in one out of four published SN Ib. Our sample of new SN does not include broad-lined SN Ic or peculiar objects. We show that this particular line profile is most probably not caused by optical depth effects, but might be due to global anisotropies in the ejecta. Prior to this work, double-peaked oxygen lines had only been reported in two peculiar SN Ib/c and interpreted as off-axis GRB-jet or unipolar blob ejections. It seems more likely that asphericities are present in a wide variety of core-collapse events and they are not strictly confined to the supernovae associated with GRB jets. Although special models have been proposed to account for the line profiles in peculiar supernovae, our investigation suggests that double-peaked profiles, and underlying disks, are not unusual. These ordinary SN from our sample have a variety of subtypes, reflecting the diversity of mass loss prior to the explosion. Indeed, most of the supernovae in our small sample that have double-peaked lines are SN IIb or SN Ib, while attempts to estimate absolute occurrence rates have to await larger samples. We note that Maeda et al. (2008) recently obtained similar results on the high frequency of double-peaked profiles in a different set of late-time spectra of stripped core-collapse SN. We recommend a meta-analysis of available nebular spectra of all core-collapse SN, spanning the full range from SN II to IIb, Ib, Ic and finally, to broad-lined SN Ic, in order to quantify the kinematics and geometry of the ejecta and to seek trends as a function of SN type. Further high-SN/N multi-epoch observations of a sample of SN II/IIb/Ib/Ic coupled with radiative transfer models should help to elucidate the observed blue- and redshifts of the line profiles.

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Fig. 1.— Selection of nebular spectra of SN IIb, SN Ib and SN Ic in their respective rest frames. SN name, type and phase of spectrum (with respect to maximum light, except for GRB 060218/SN 2006aj which is referenced to time of GRB burst) are marked. For clarity we show only the latest spectrum of each SN (except for SN 2004ao), even though we obtained several for some SN. Also the main nebular emission lines of $[\text{Mg I} \lambda 4571$, $[\text{O I}] \lambda\lambda 6300, 6363$, $[\text{Ca II}] \lambda\lambda 7291, 7324$ and $\text{OI} \lambda 7774$ are marked at the very top.
Fig. 2.— Montage of SN with double-peaked oxygen profiles in velocity space. SN name, type and phase of spectrum (with respect to maximum light). SN 1984L, a SN Ib is from Schlegel & Kirshner (1989). SN 1994I (Filippenko et al. 1995) that exhibits a simple parabolic oxygen line profile is plotted for comparison at the bottom. The dashed line marks zero velocity with respect to 6300 Å while the dotted line corresponds to [O I] λ6363. For SN 2004ao, we plot the scaled profiles of O I λ7774 and [Mg I] λ4571 (blue lines), single-transition lines that also exhibit the two peaks. As discussed in the text (§ 3.1), the two horns are unlikely to be due to the doublet nature of [O I] λ6300,6363. Note the large blueshift in SN 2005bf compared to the other SN.
Fig. 3.— Evolution of the intensity ratio of the blue to the red horn (\( I(\text{Blue Horn})/I(\text{Red Horn}) \)) as a function of time for SN 2004ao (filled circles). The intensity ratio decreases over time. In contrast, if the blue and red horns corresponded to [O I] \( \lambda 6300 \) and [O I] \( \lambda 6363 \) the intensity ratio would increase over time, as observed in SN 1987A (empty circles, from Li & McCray 1992) and 1988A (Spyromilio 1991).